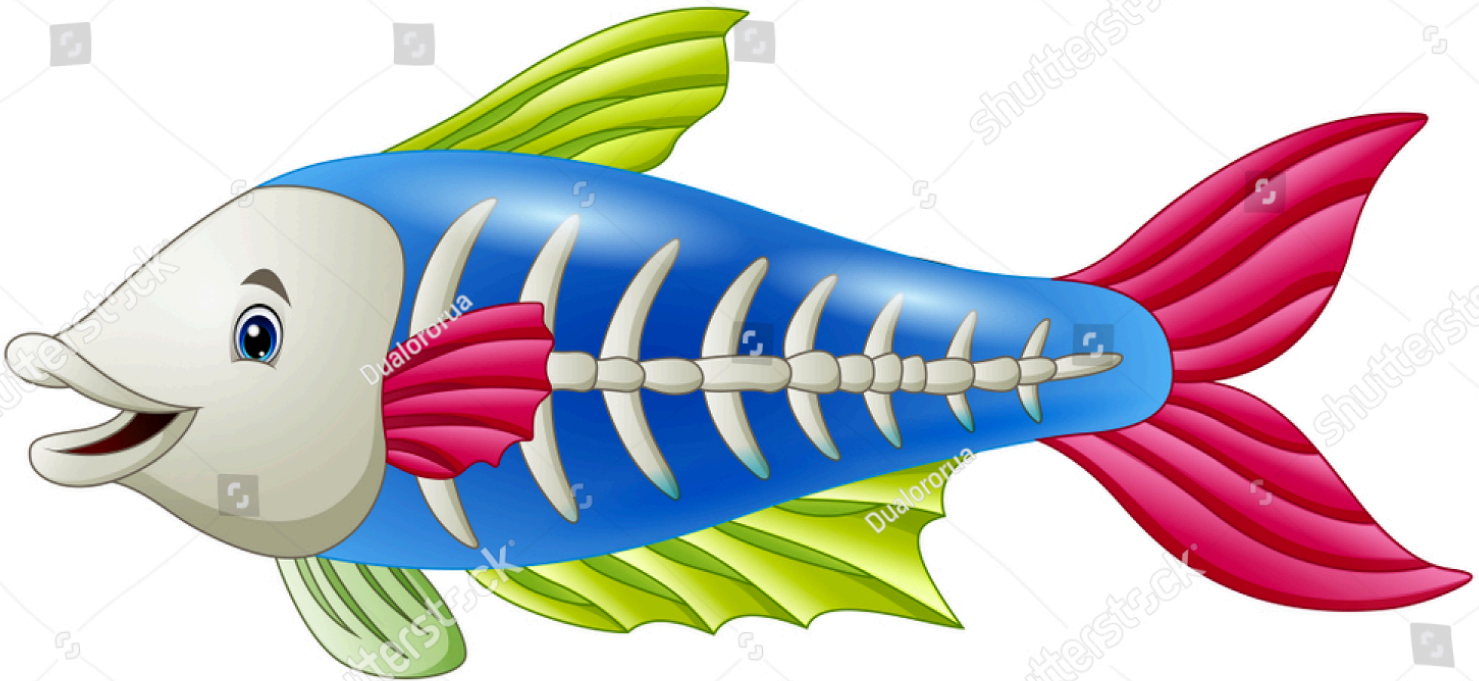




Atomic Physics

Chapter 6 X ray

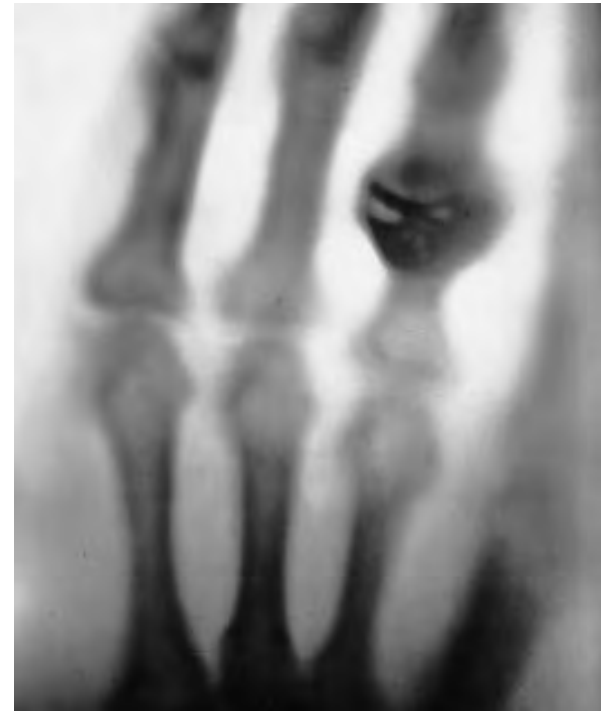


6.1 The discovery of X ray

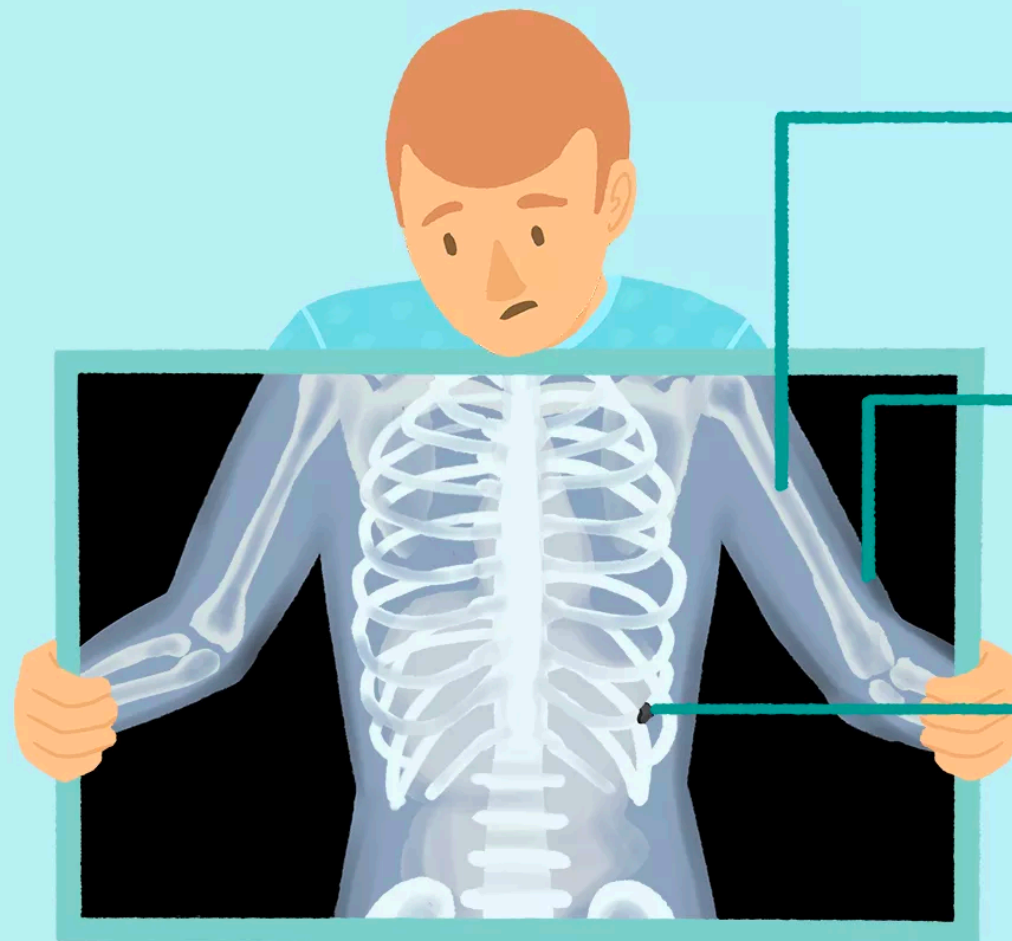


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X-rays were discovered in 1895 by the German physicist Wilhelm Roentgen. He found that a beam of high-speed electrons striking a metal target produced a new and extremely penetrating type of radiation.



How to Read an X-Ray



Bones

Block radiation,
appear white

Muscle, Fluid, Fats

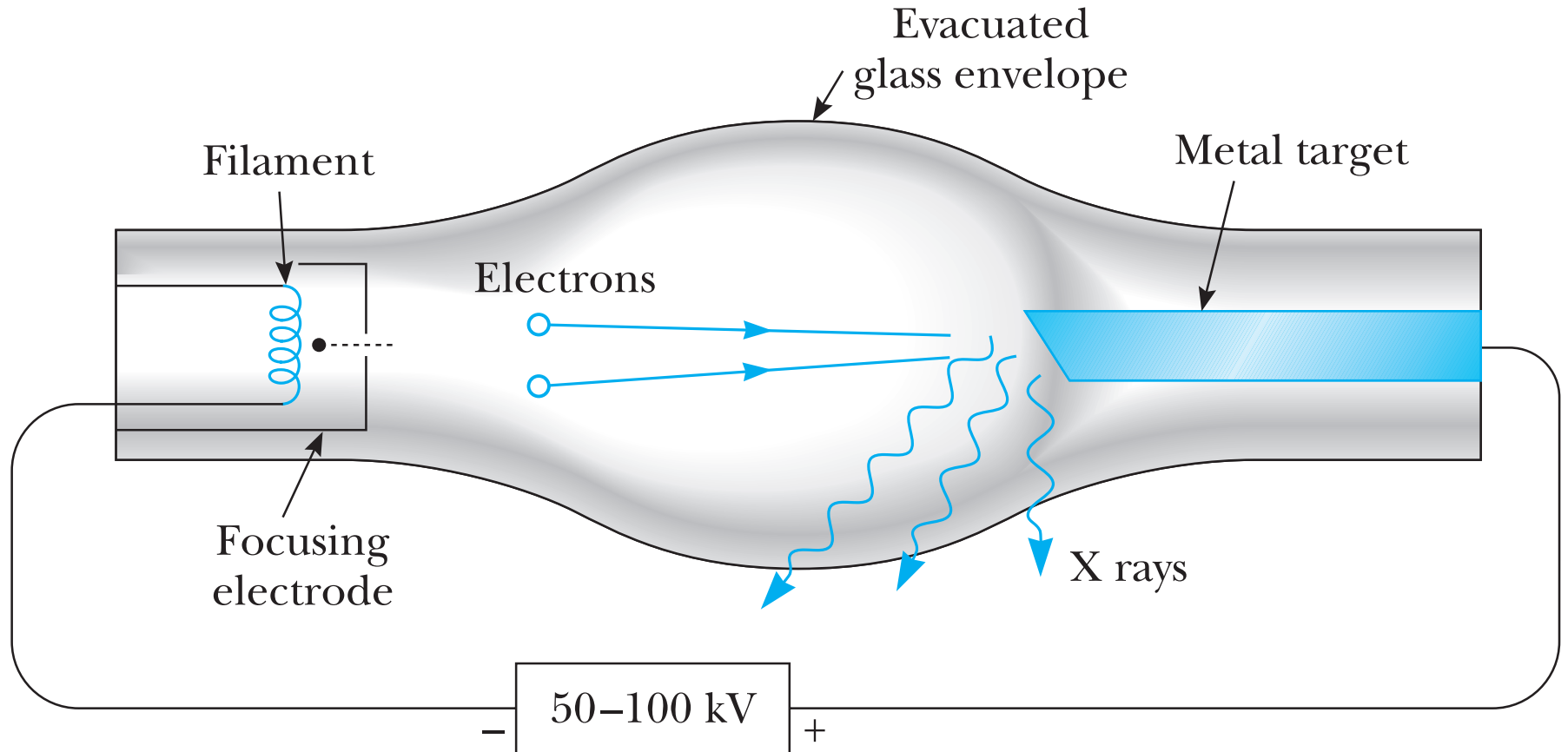
Some radiation passes
through, appear grey

Breaks, Fractures

Radiation passes
completely through,
appear black

6.1 The discovery of X ray

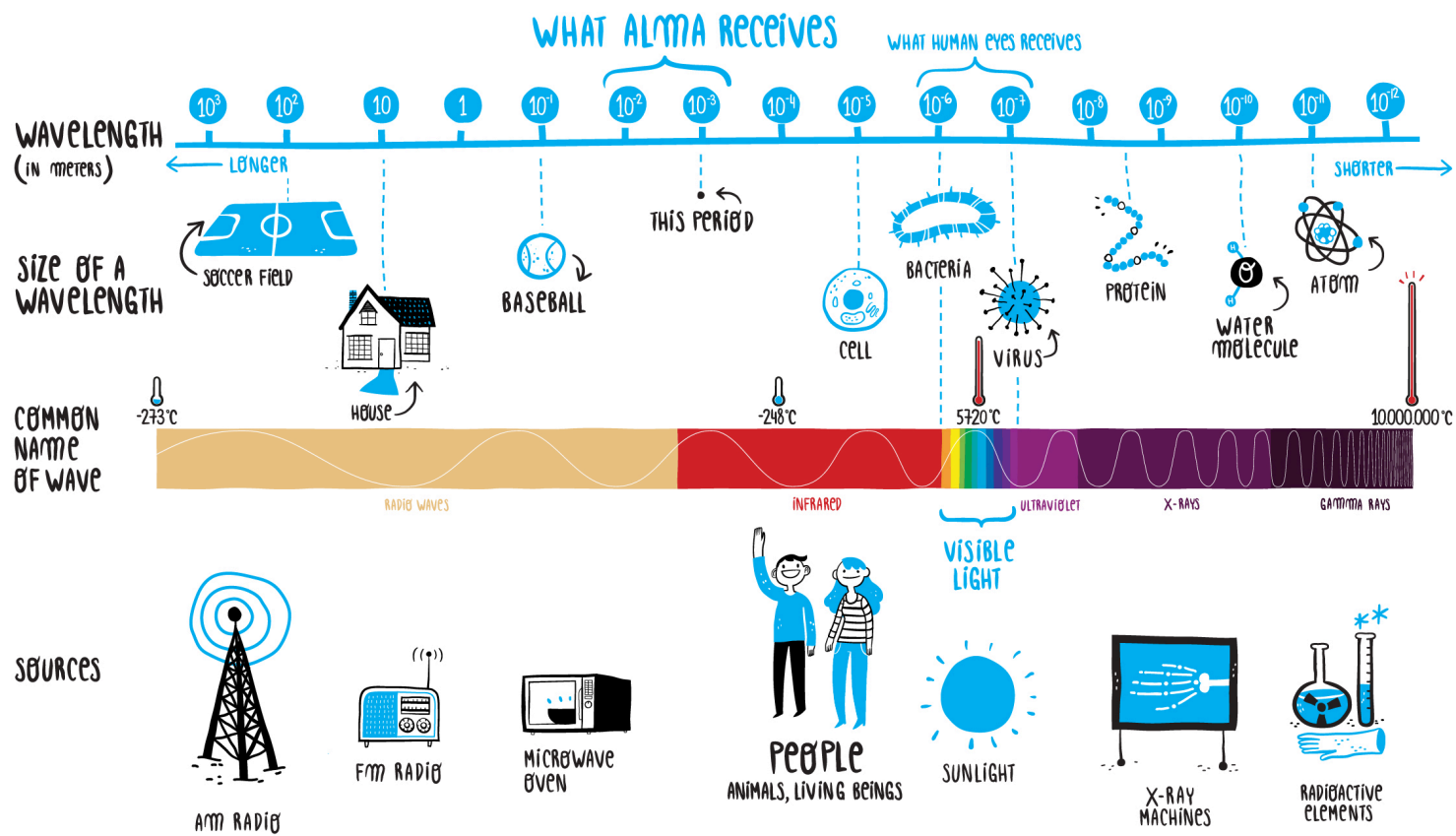
X-rays Tube



6.1 The discovery of X ray

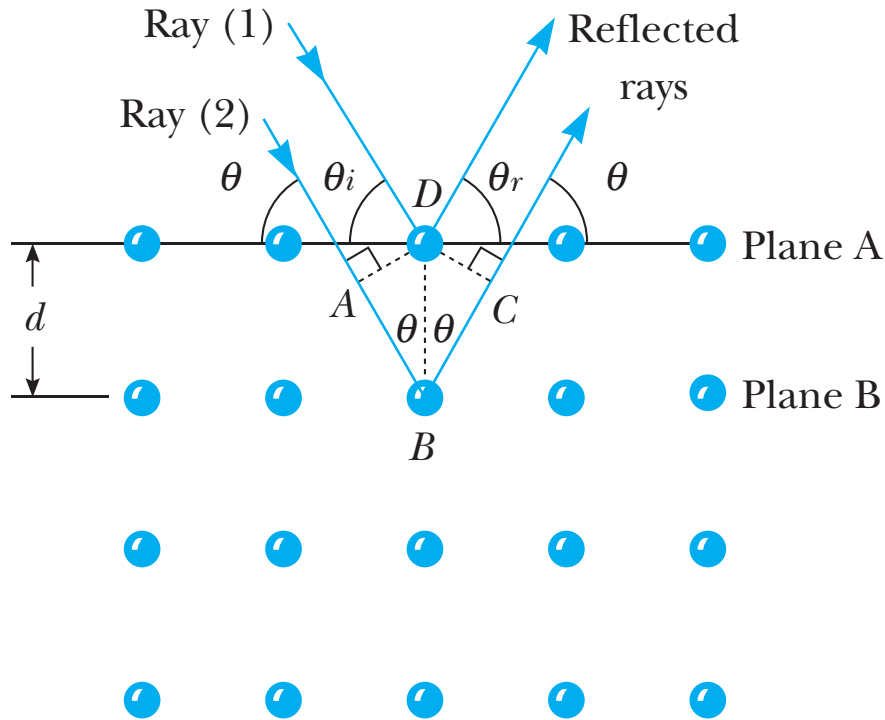
Electromagnetic wave

THE ELECTROMAGNETIC SPECTRUM



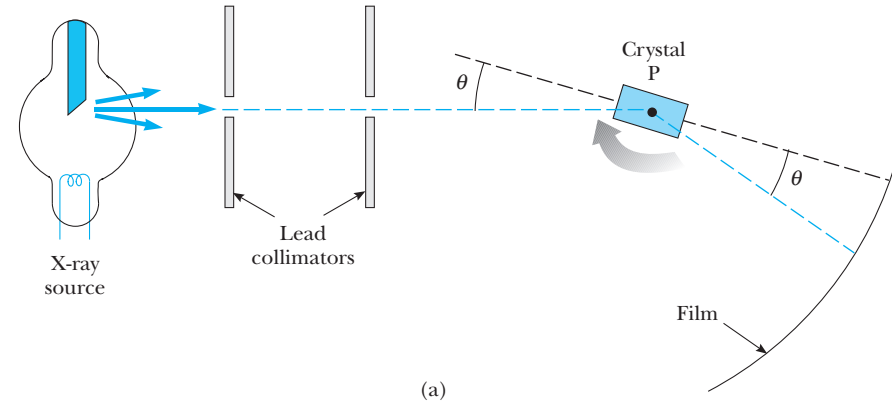
6.1 The discovery of X ray

X ray diffraction



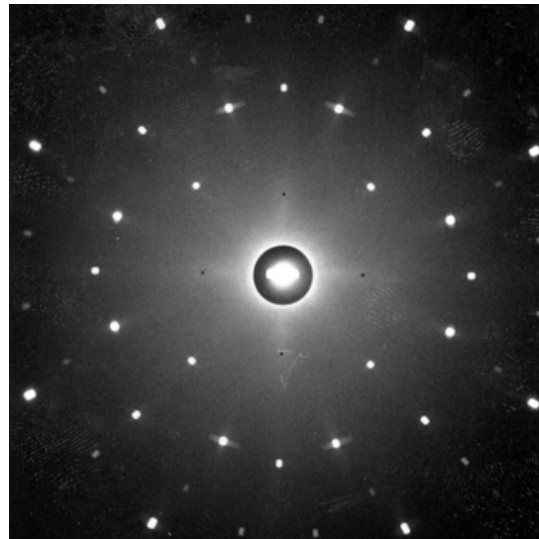
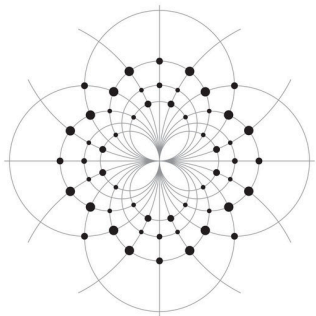
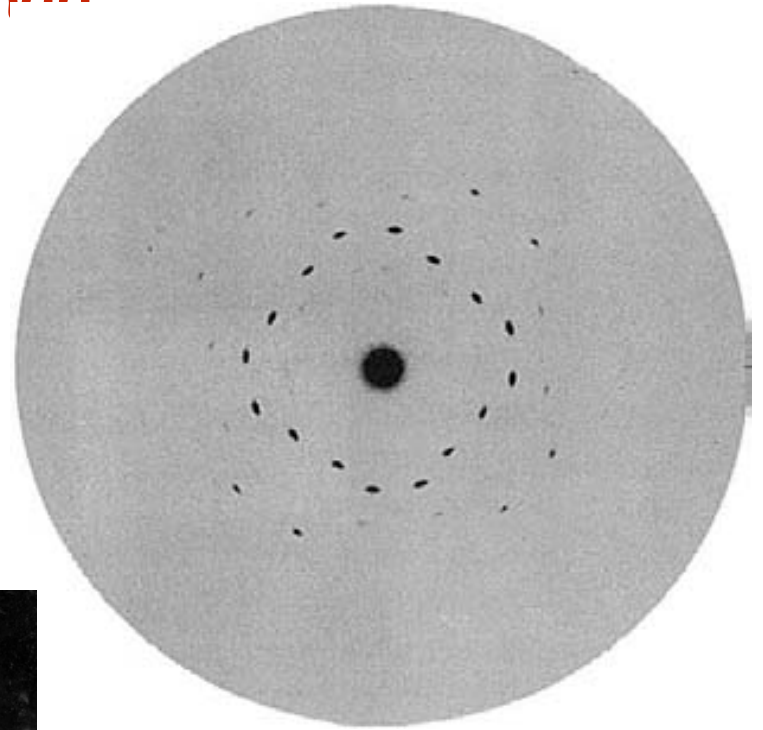
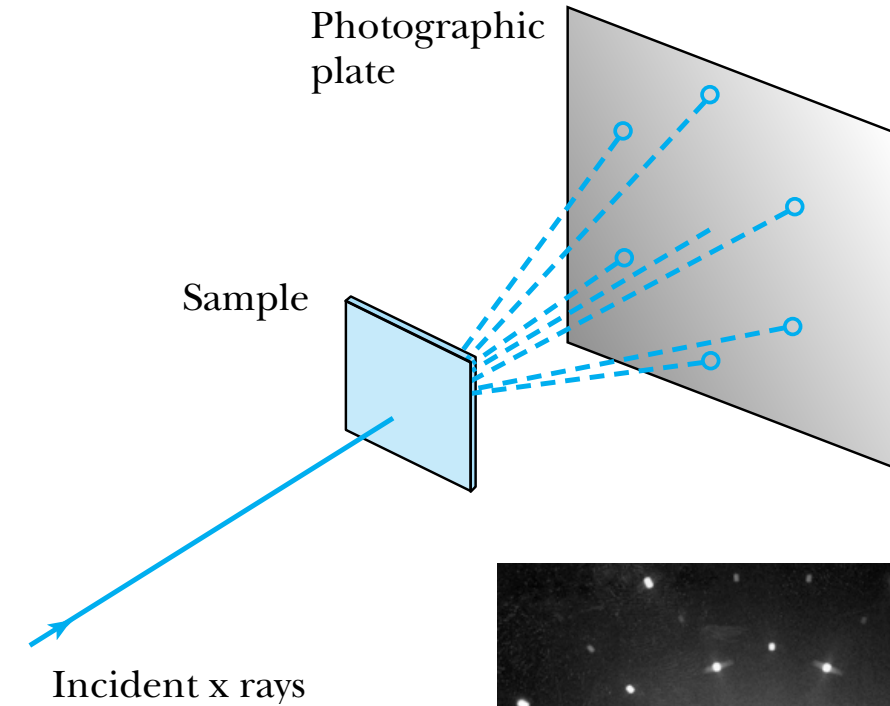
Bragg equation

$$n\lambda = 2d \sin \theta \quad n = 1, 2, 3, \dots$$



6.1 The discovery of X ray

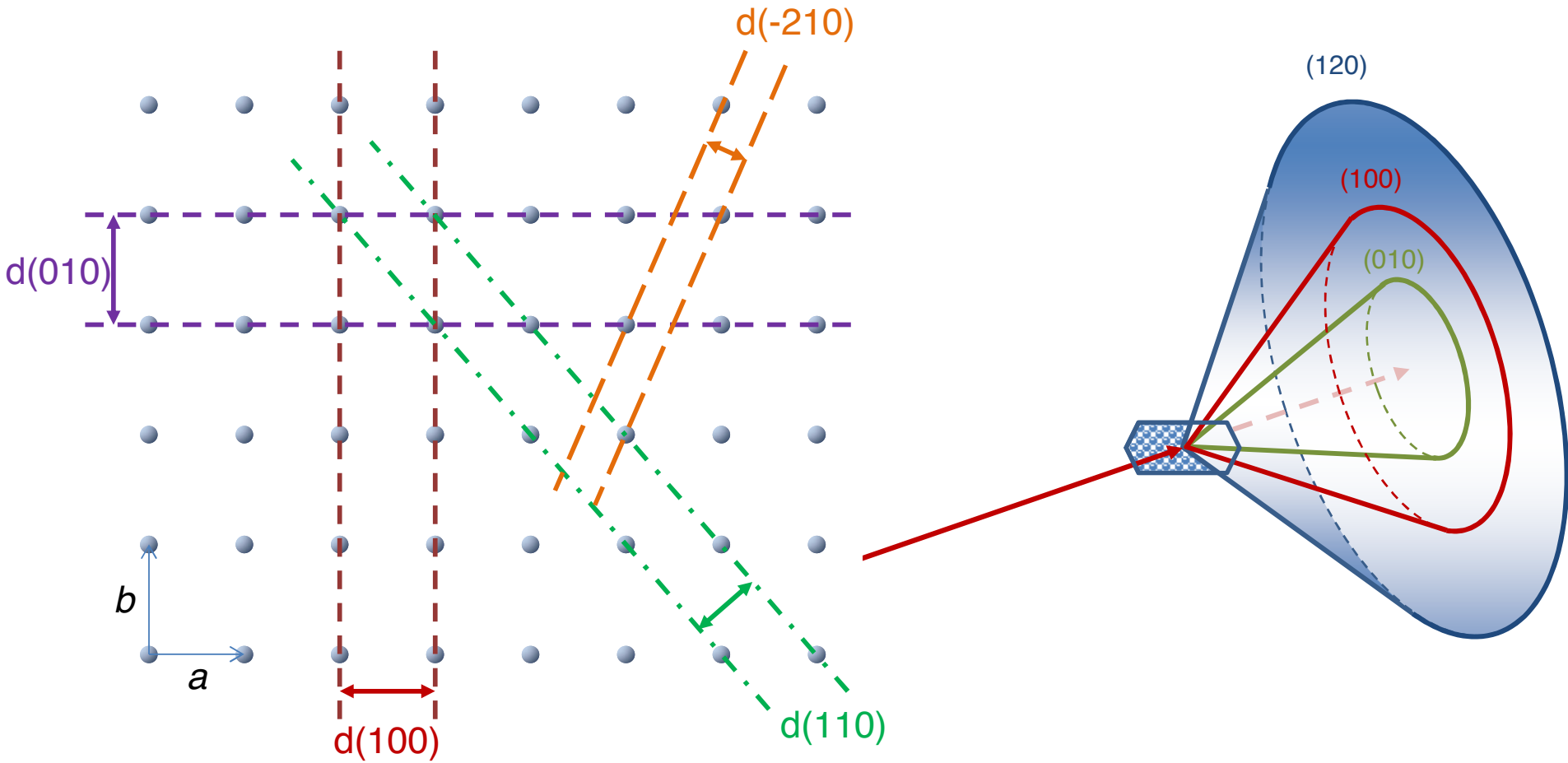
Laue diffraction transmission method



The diffraction of X ray



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6.1 The discovery of X ray

The polarization of X ray





The polarization of X ray



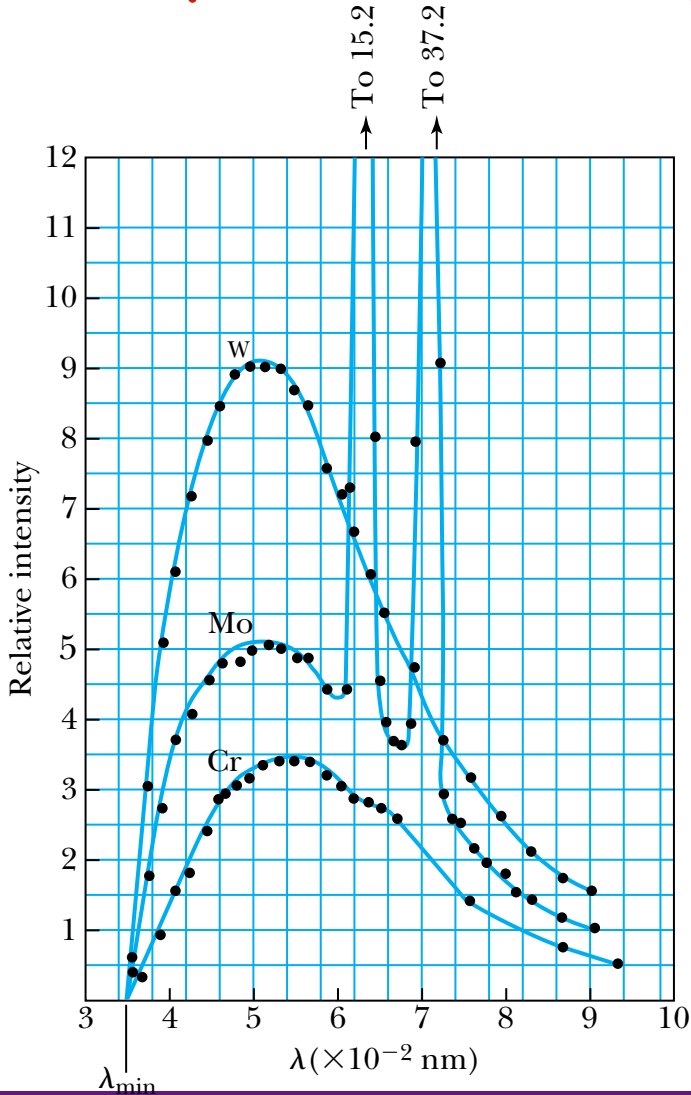
POLARISATION

6.2 X ray production

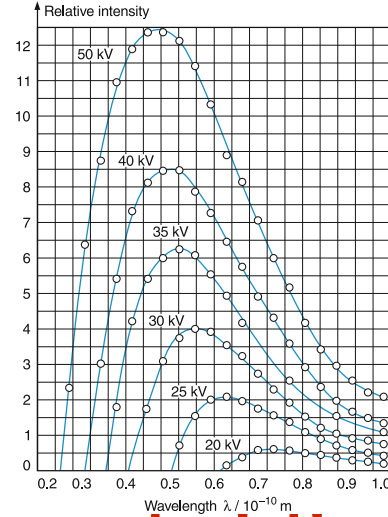


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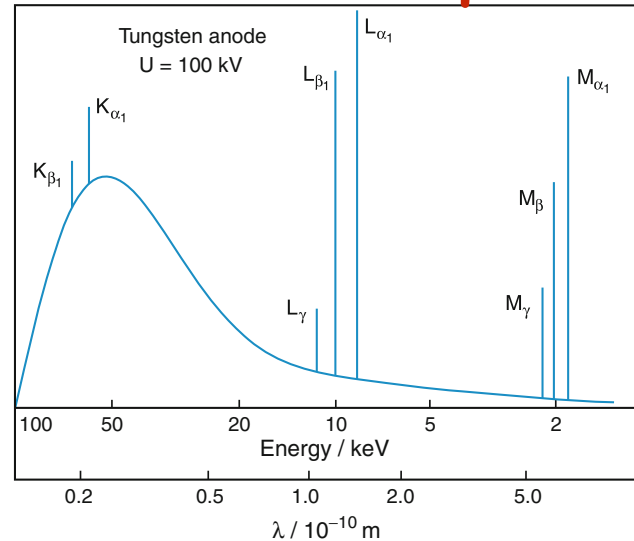
The spectrum of X ray



The continuous spectrum



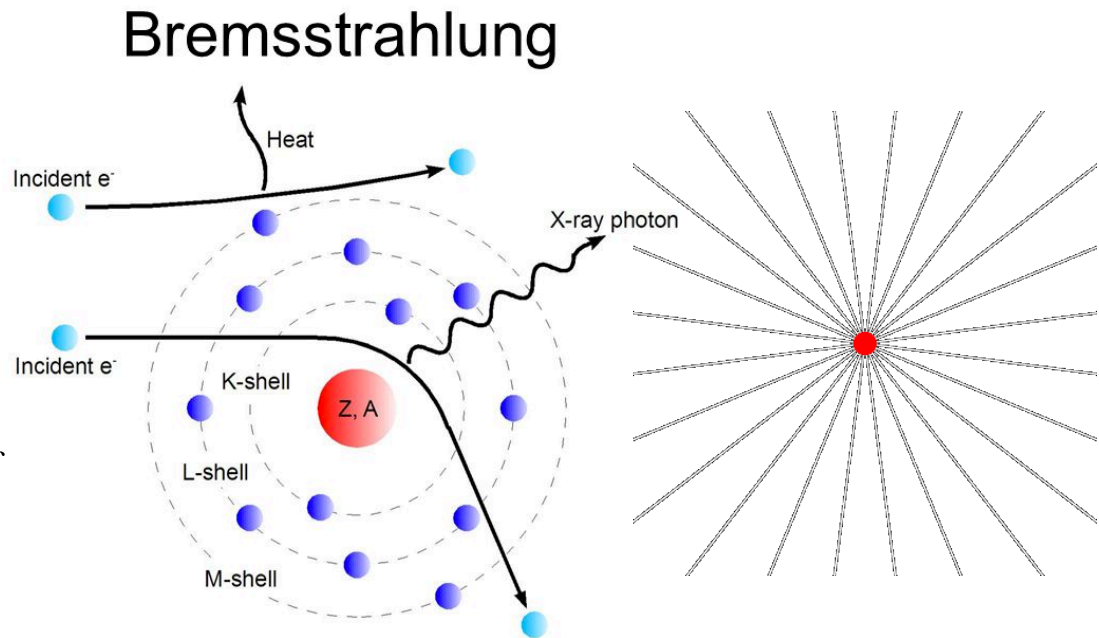
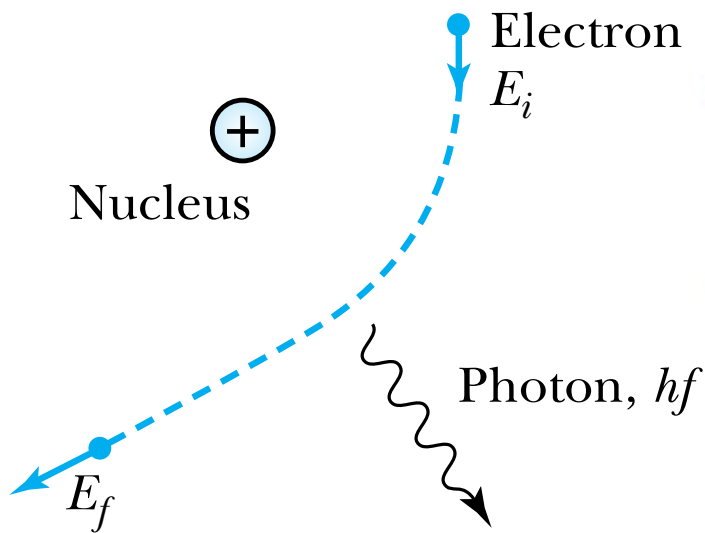
The characteristic spectrum

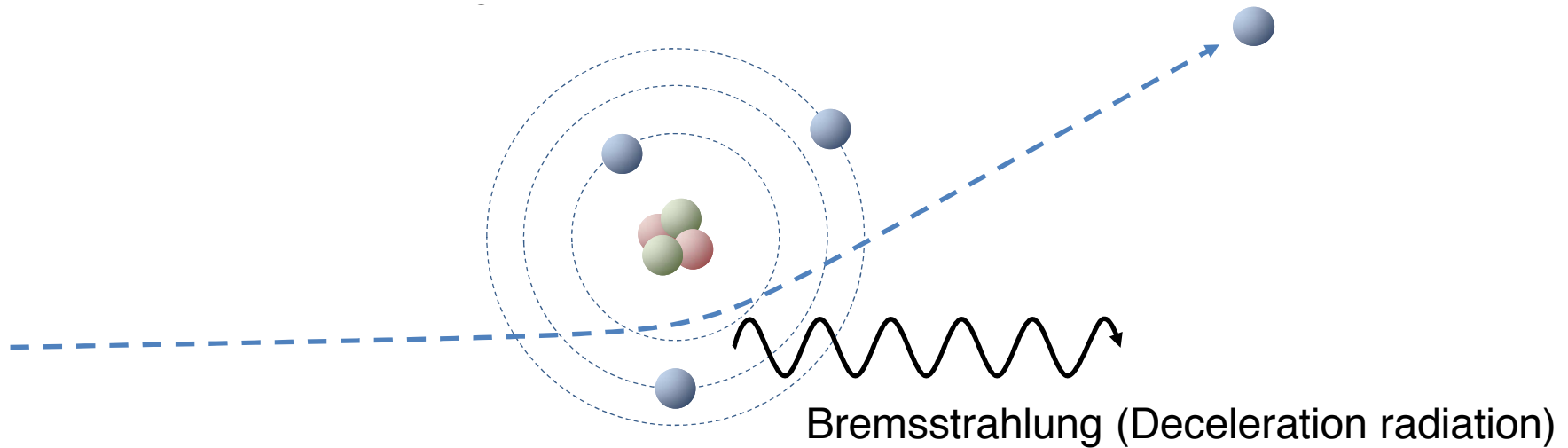


6.2 X ray production

The continuous spectrum:

An energetic electron passing through matter will radiate photons and lose kinetic energy. The process by which photons are emitted by an electron slowing down is called **bremsstrahlung**, from the German word for “braking radiation.”





Electron is deflected and decelerated by the atomic nucleus. (Inelastic scattering).

Deflected electron emits electromagnetic radiation. Wavelength depends on the loss of energy.

6.2 X ray production



The minimum wavelength is due to the **inverse photoelectric effect**. The conservation of energy requires that the electron kinetic energy equal the maximum photon energy:

$$eV_0 = hf_{\max} = \frac{hc}{\lambda_{\min}}$$

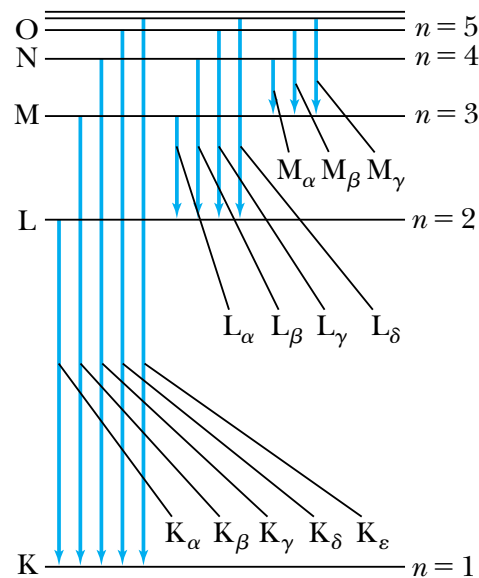
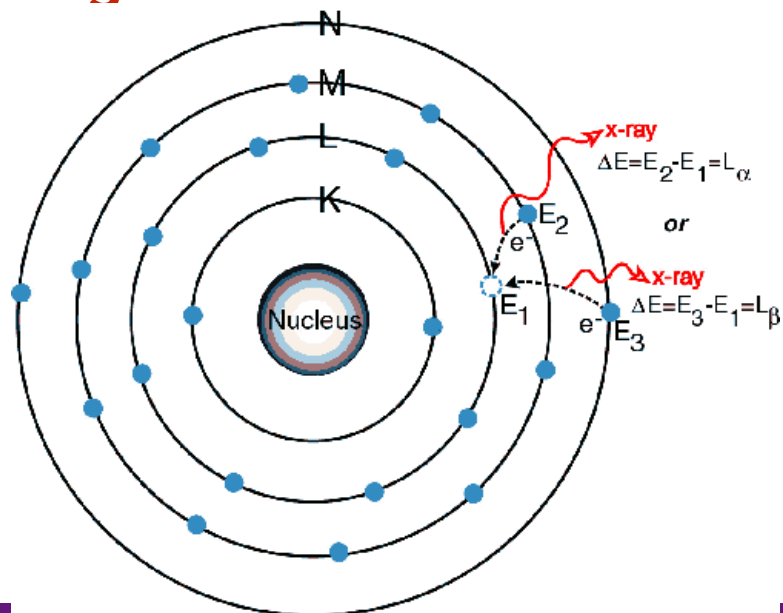
Therefore, the minimum wavelength (Duane-Hunt rule) is

$$\lambda_{\min} = \frac{hc}{e} \frac{1}{V_0} = \frac{1.240 \times 10^{-6} \text{ V} \cdot \text{m}}{V_0}$$

6.2 X ray production

The characteristic spectrum:

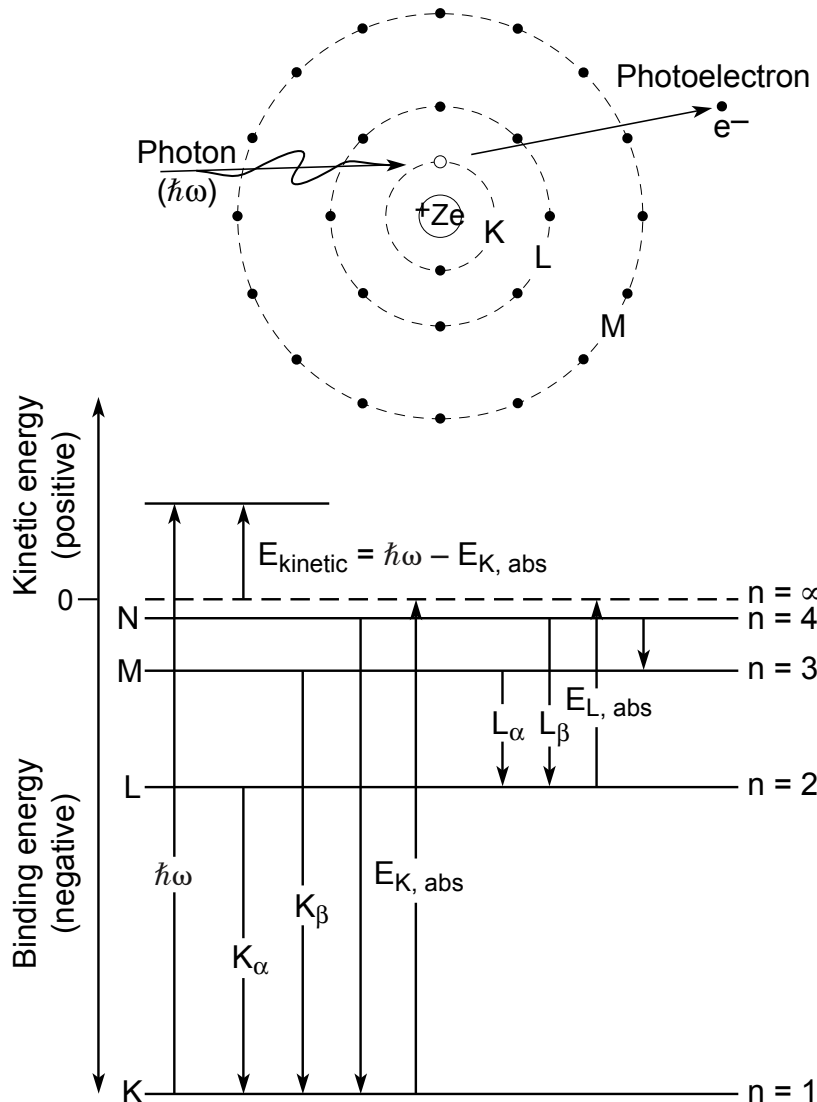
The atom is most stable in its lowest energy state or ground state, so it is likely that an electron from one of the higher shells will change its state and fill the inner-shell vacancy at lower energy, emitting radiation as it changes its state.



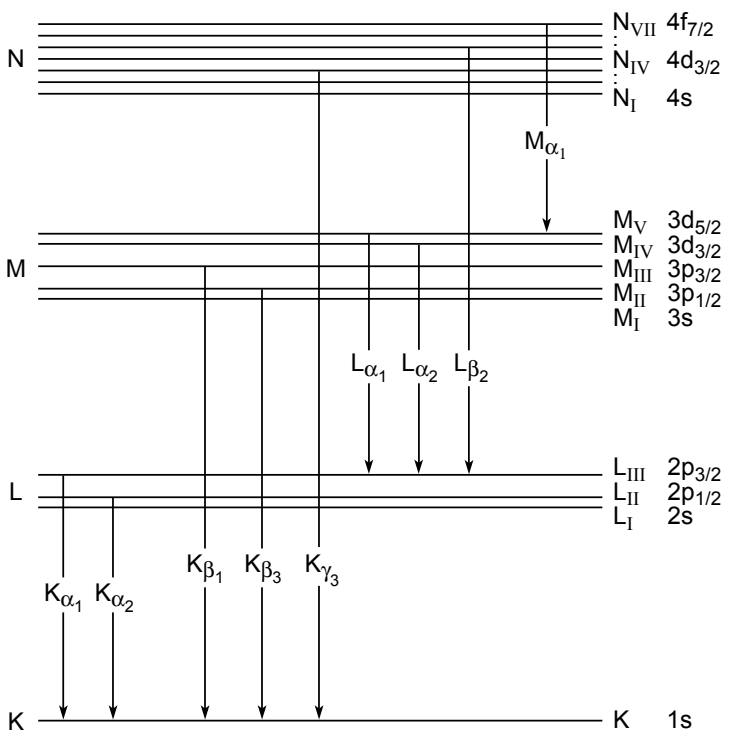
The characteristic spectrum



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n	l	j
4	3	7/2
4	3	5/2
...
4	0	1/2
3	2	5/2
3	2	3/2
3	1	3/2
3	1	1/2
3	0	1/2
2	1	3/2
2	1	1/2
2	0	1/2
1	0	1/2



Absorption edges for copper (Z = 29):

- $E_{N_1, \text{abs}} = 7.7 \text{ eV}$
- $E_{M_3, \text{abs}} = 75 \text{ eV}$
- $E_{M_1, \text{abs}} = 123 \text{ eV}$
- $E_{L_3, \text{abs}} = 933 \text{ eV}$
- $E_{L_2, \text{abs}} = 952 \text{ eV}$
- $E_{L_1, \text{abs}} = 1,097 \text{ eV}$
- $E_{K, \text{abs}} = 8,979 \text{ eV} (1.381\text{\AA})$

- Cu $K_{\alpha_1} = 8,048 \text{ eV} (1.541\text{\AA})$
- Cu $K_{\alpha_2} = 8,028 \text{ eV} (1.544\text{\AA})$
- Cu $K_{\beta_1} = 8,905 \text{ eV}$
- Cu $L_{\alpha_1} = 930 \text{ eV}$
- Cu $L_{\alpha_2} = 930 \text{ eV}$
- Cu $L_{\beta_1} = 950 \text{ eV}$

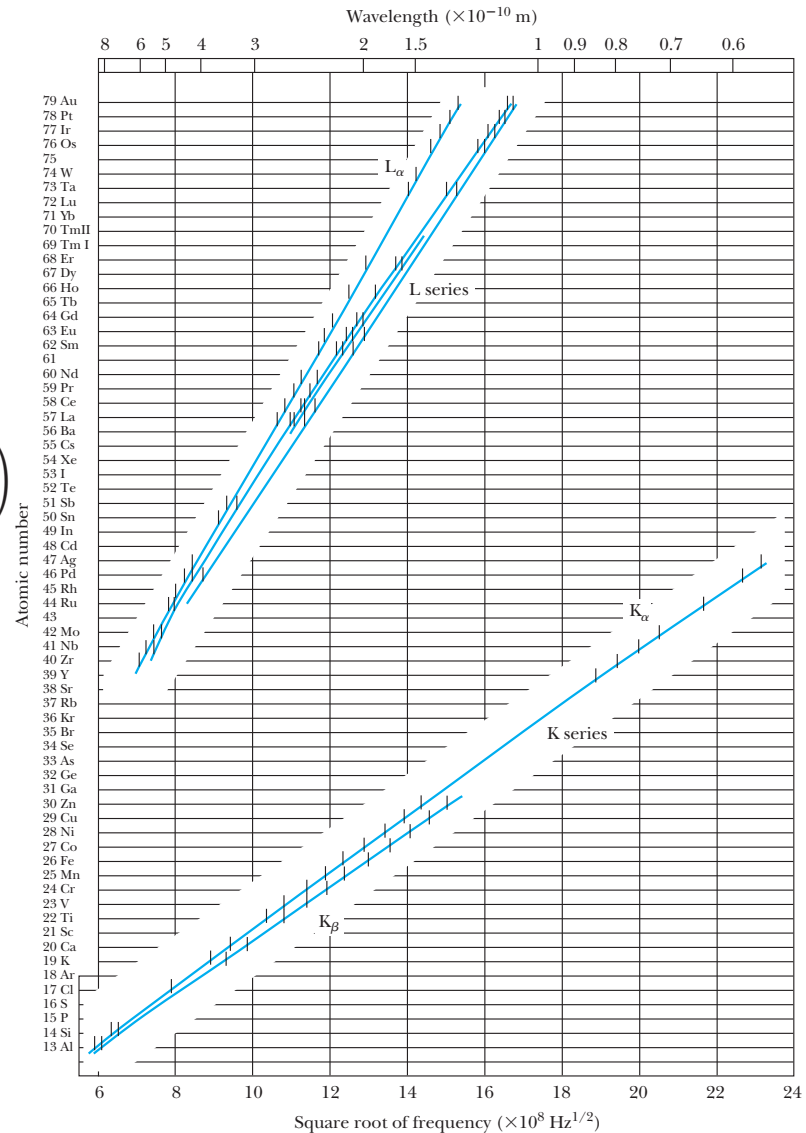
6.2 X ray production

Moseley formula

$$\frac{1}{\lambda_{K\alpha}} = R(Z - 1)^2 \left(\frac{1}{1^2} - \frac{1}{2^2} \right) = \frac{3}{4} R(Z - 1)^2$$

$$f_{K\alpha} = \frac{c}{\lambda_{K\alpha}} = \frac{3cR}{4} (Z - 1)^2$$

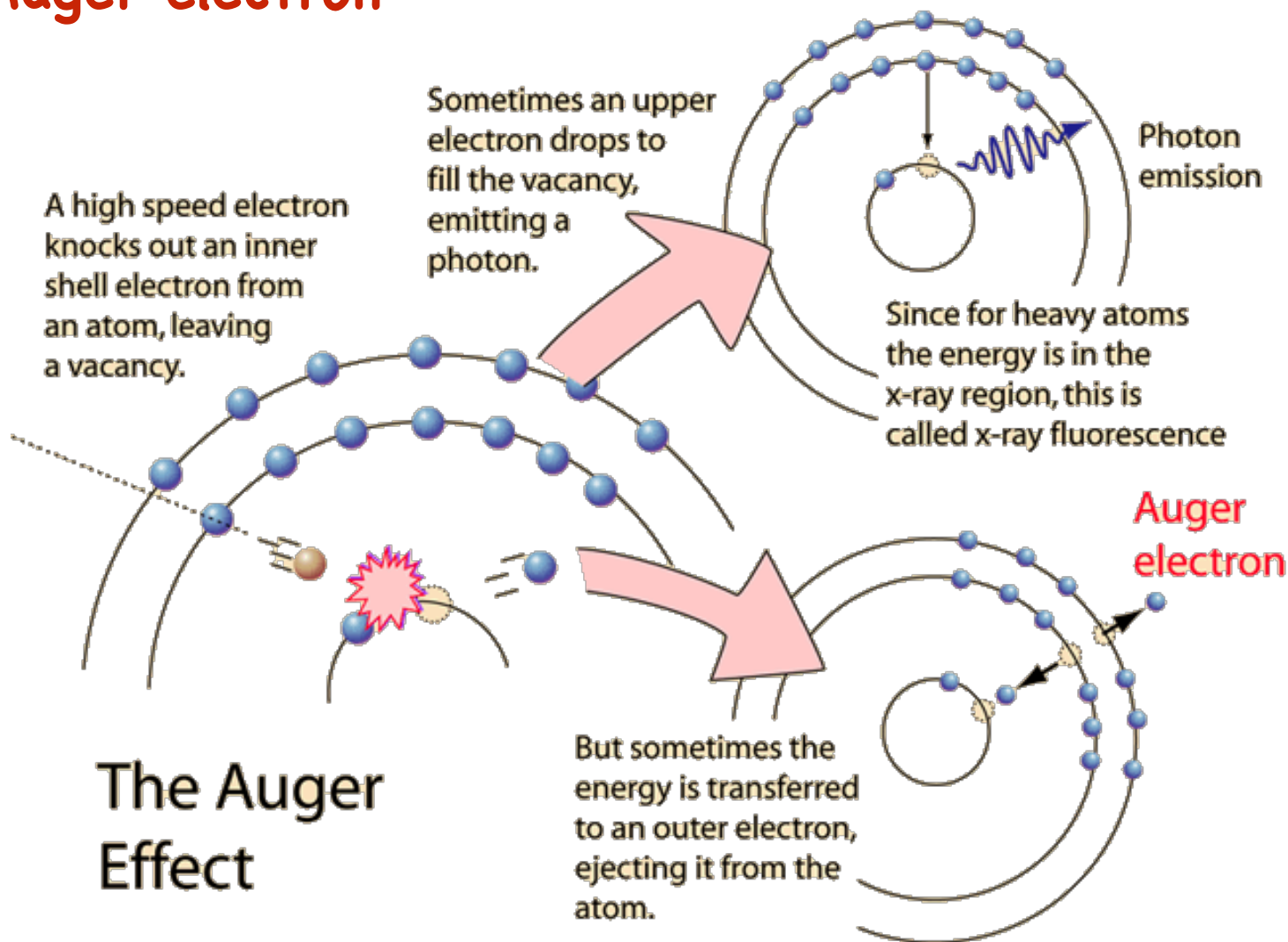
$$\frac{1}{\lambda_K} = R(Z - 1)^2 \left(\frac{1}{1^2} - \frac{1}{n^2} \right) = R(Z - 1)^2 \left(1 - \frac{1}{n^2} \right)$$



6.2 X ray production



Auger electron

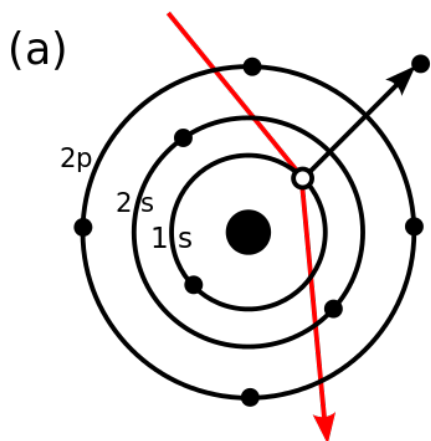


6.2 X ray production

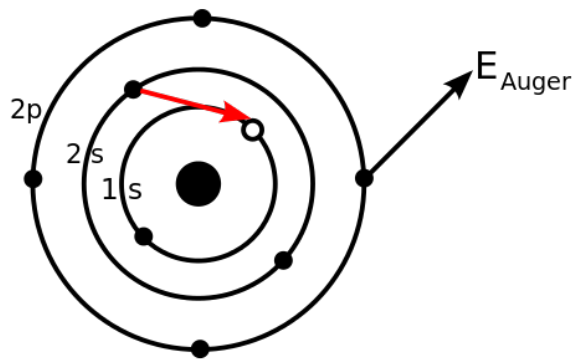


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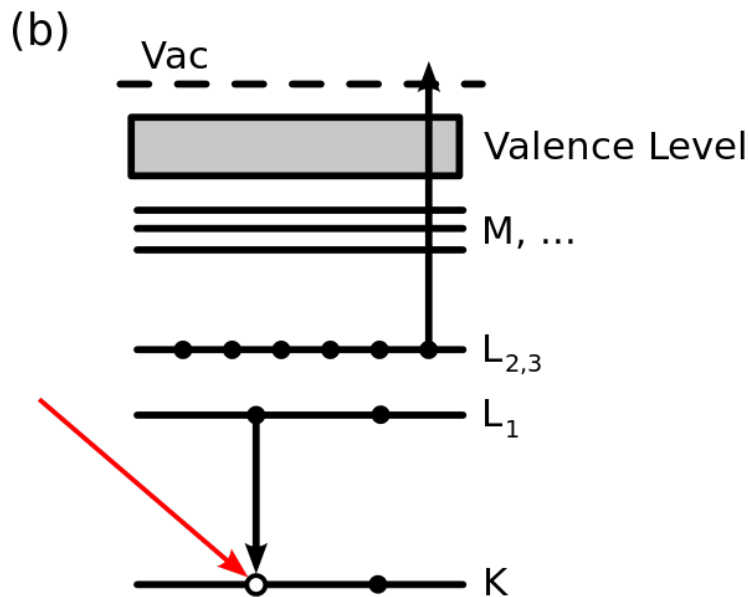
Auger electron



Electron collision



Auger electron emission

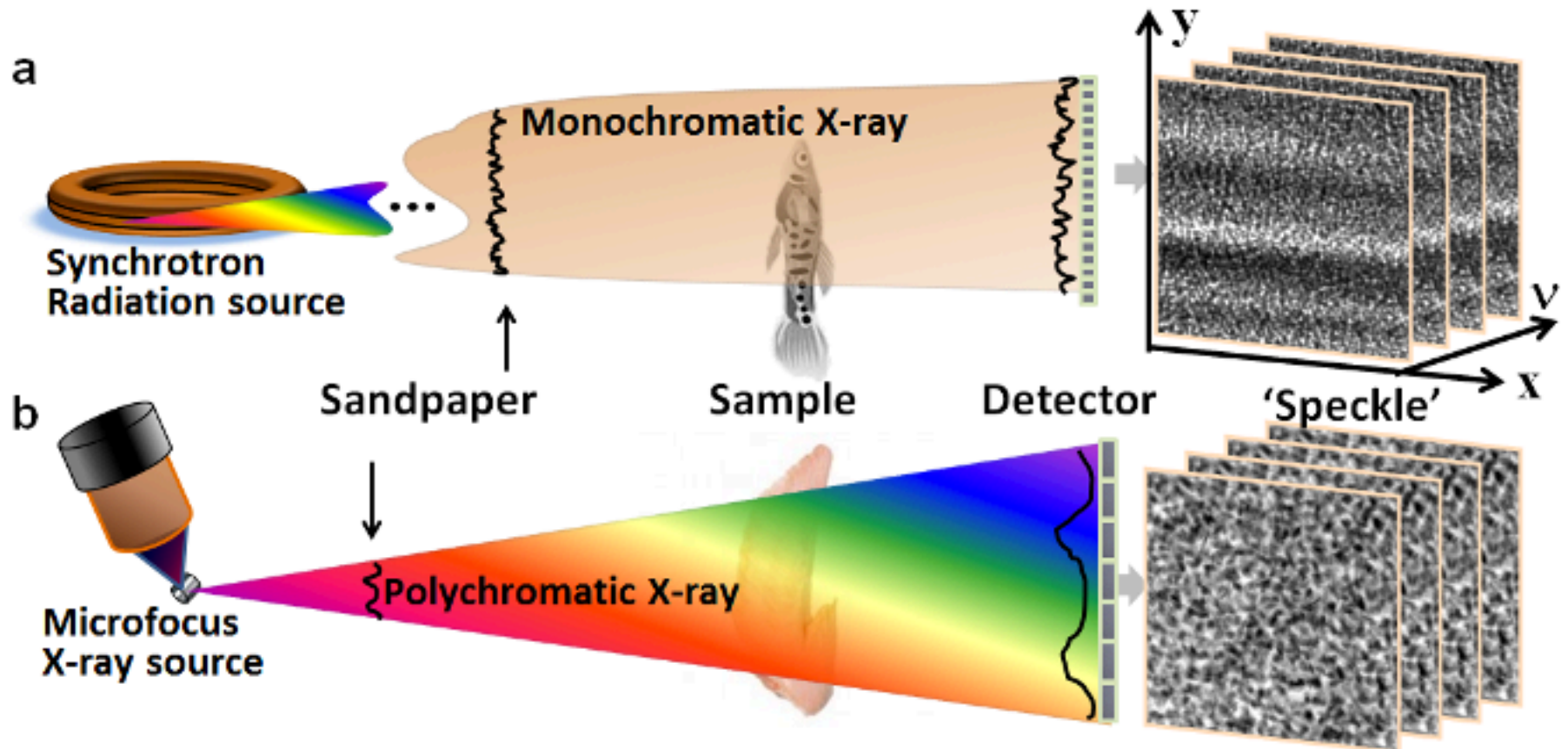


The kinetic energy of Auger electron

$$E_{ae} = E_K - E_{L_1} - E_{L_{2,3}}$$

6.2 X ray production

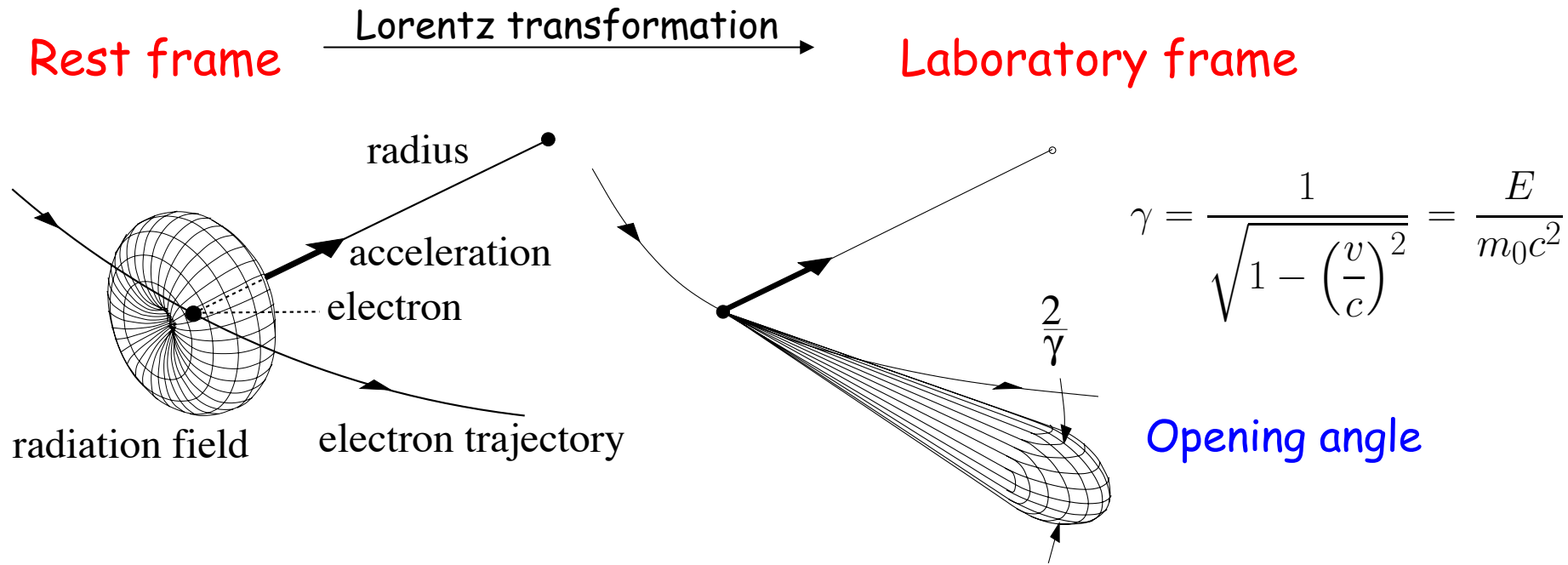
Synchrotron Radiation: In cyclic accelerators, when charged particles are accelerated, they radiate electromagnetic energy called synchrotron radiation.



Synchrotron Radiation



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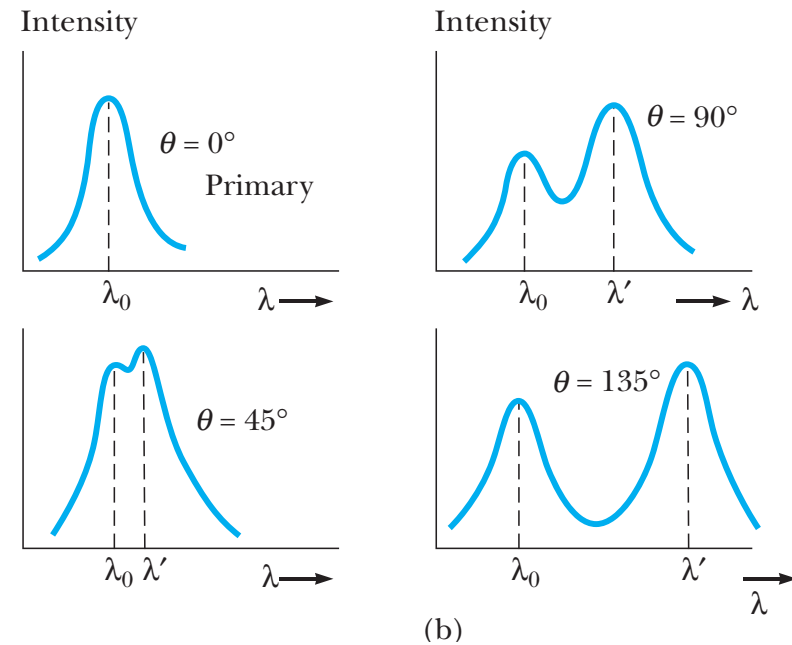
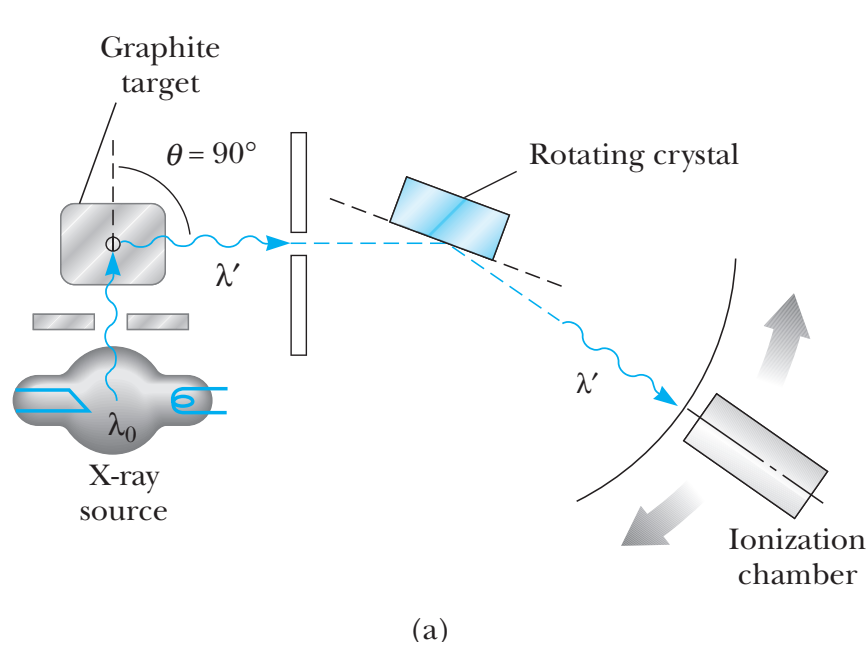
The very high intensity of the source yields images with a high signal-to-noise ratio on short time-scales, which enables fast radiographic investigations.

The beam can be easily monochromated. This allows correlations between attenuation values and the chemical constituents of the sample

The option to vary the energy of the radiation enables the investigation of objects with very different absorption coefficients within the same measuring environment.

6.3 Compton scattering

At backward-scattering angles, there appeared to be a component of the emitted radiation (called a modified wave) that had a longer wavelength than the original primary (unmodified) wave.



6.3 Compton scattering



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Photon Interactions

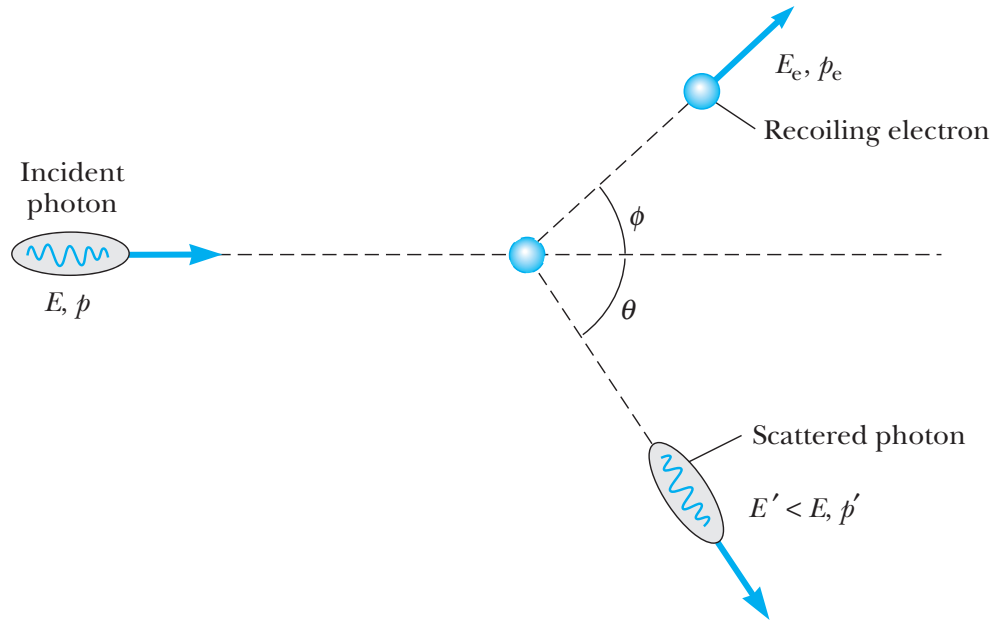


6.3 Compton scattering



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Diagram representing Compton scattering of a photon by an electron.



Conservation of energy

$$E + m_e c^2 = E' + E_e$$

Conservation of momentum

$$p = p' \cos \theta + p_e \cos \phi$$

$$p' \sin \theta = p_e \sin \phi$$

6.3 Compton scattering



Therefore

$$p_e^2 = (p')^2 + p^2 - 2pp' \cos \theta$$

With De Broglie relation

$$p_{\text{photon}} = \frac{E}{c} = \frac{hf}{c} = \frac{h}{\lambda}$$

We have

$$E_e = hf - hf' + m_e c^2$$

$$E_e^2 = p_e^2 c^2 + m_e^2 c^4$$

$$p_e^2 = \left(\frac{hf'}{c} \right)^2 + \left(\frac{hf}{c} \right)^2 - \frac{2h^2 ff'}{c^2} \cos \theta$$

Finally

$$\lambda' - \lambda_0 = \boxed{\frac{h}{m_e c}} (1 - \cos \theta) \rightarrow \text{Compton wavelength}$$

$\lambda_C = 2.426 \times 10^{-3} \text{ nm}$

6.3 Compton scattering

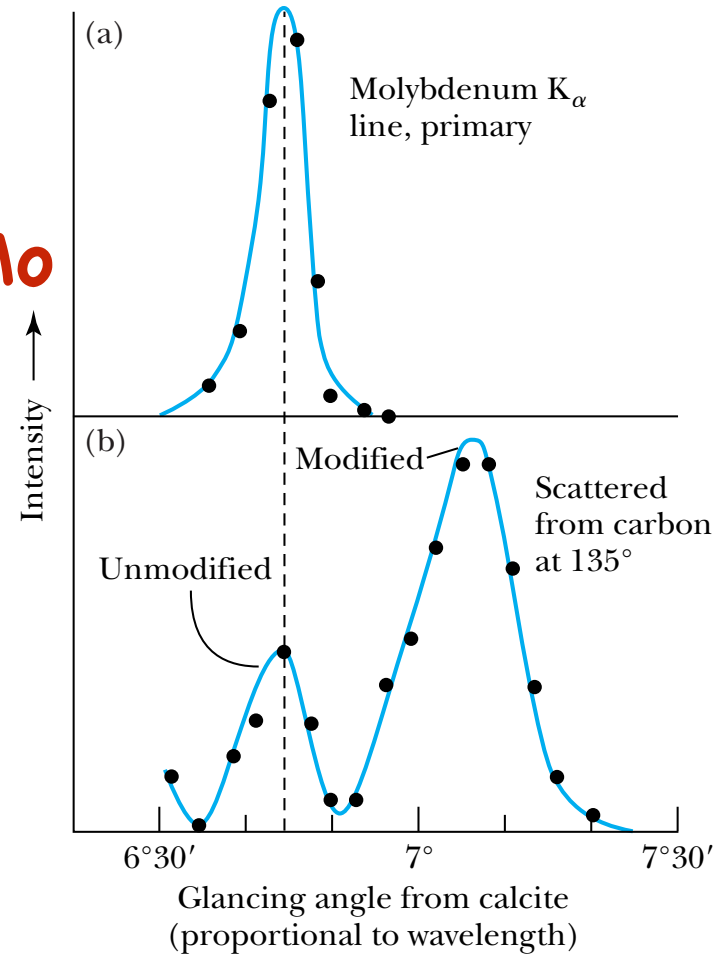


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Compton's original data showing

1. the primary x-ray beam from Mo unscattered

2. the scattered spectrum from carbon at 135° showing both the modified and unmodified wave.



In the photoelectric effect, bremsstrahlung, and the Compton effect, we have studied exchanges of energy between photons and electrons. Have we covered all possible mechanisms?

For example, can the kinetic energy of a photon be converted into particle mass and vice versa?

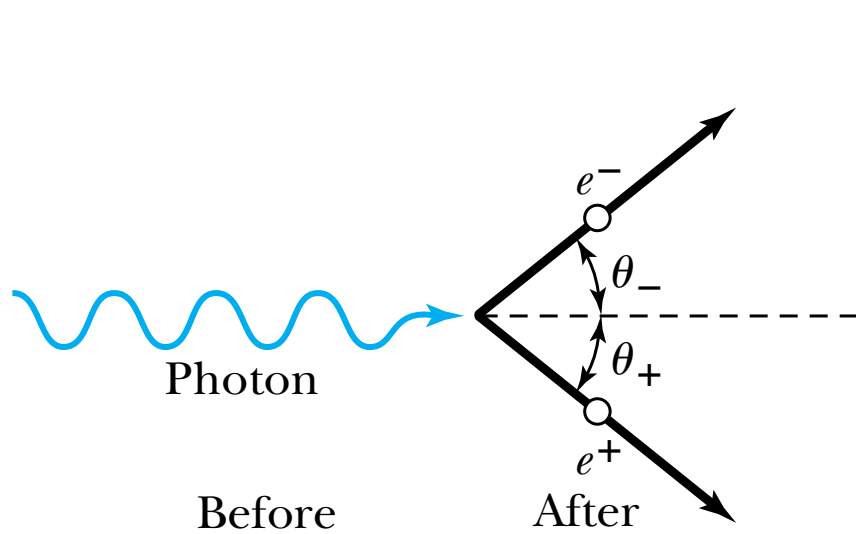
It would appear that if none of the conservation laws are violated, then such a process should be possible.

Experiments show that a photon's energy can be converted entirely into an electron and a positron in a process called pair production. The reaction is

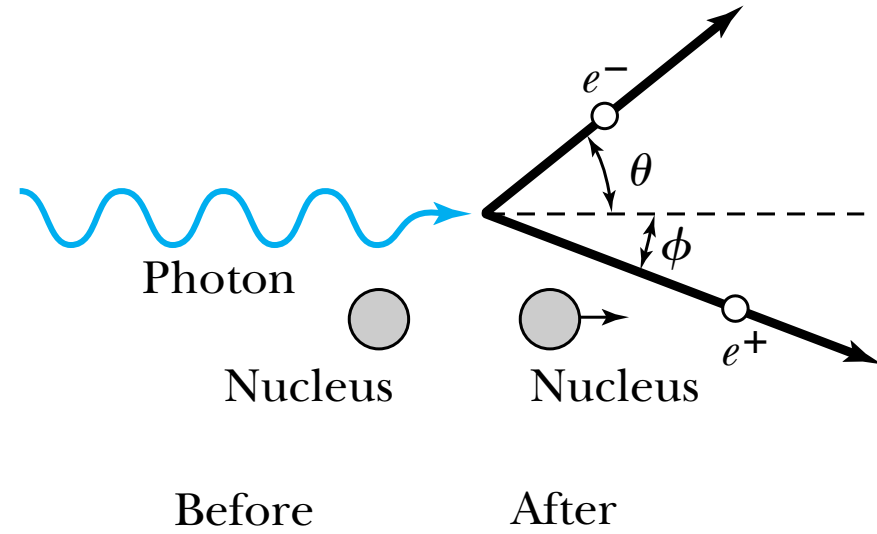


However, this process occurs only when the photon passes through matter, because energy and momentum would not be conserved if the reaction took place in isolation. The missing momentum must be supplied by interaction with a nearby massive object such as a nucleus.

Pair Production



(a) Free space (**cannot occur**)



(b) Beside nucleus

Consider the conversion of a photon into an electron and a positron that takes place inside an atom where the electric field of a nucleus is large. The nucleus recoils and takes away a negligible amount of energy but a considerable amount of momentum. The conservation of energy will now be

$$hf = E_+ + E_- + \text{K.E. (nucleus)}$$

The photon energy must be at least equal to $2m_e c^2$ in order to create the masses of the electron and positron.

$$hf > 2m_e c^2 = 1.022 \text{ MeV} \quad (\text{for pair production})$$

Pair Annihilation

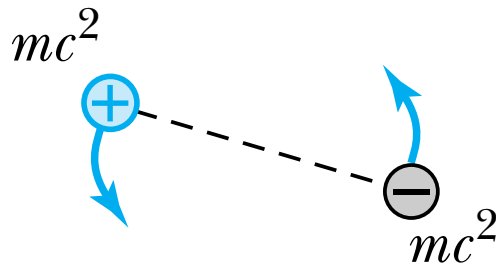


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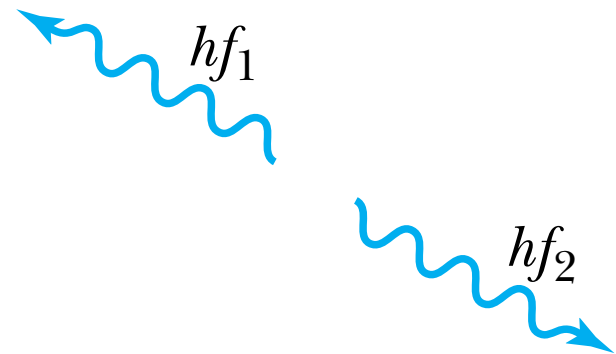
Eventually the electron and positron annihilate each other (typically in 10^{-10} s), producing electromagnetic radiation (photons). The process



is called **pair annihilation**.



(a) Positronium,
before decay
(schematic only)



(b) After annihilation

The conservation laws for the process will be

$$2m_e c^2 \approx hf_1 + hf_2$$

$$0 = \frac{hf_1}{c} - \frac{hf_2}{c}$$

The frequencies are identical

$$hf = m_e c^2 = 0.511 \text{ MeV}$$

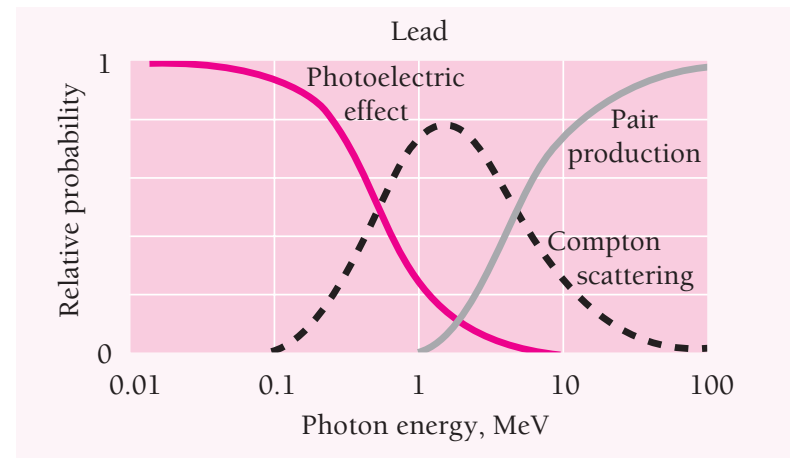
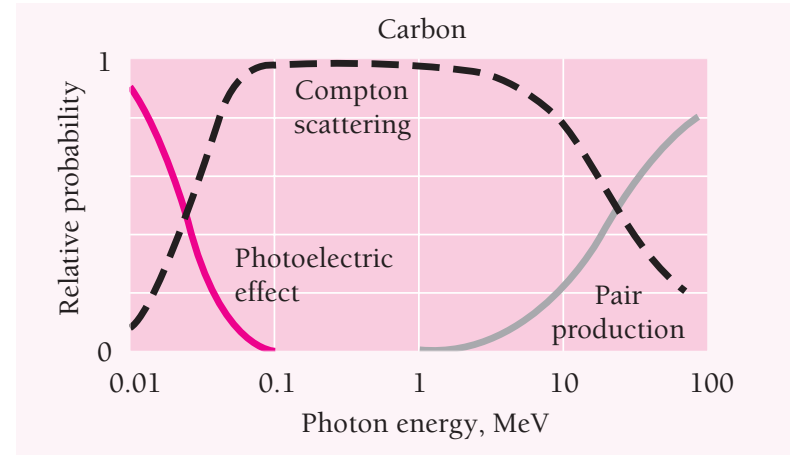
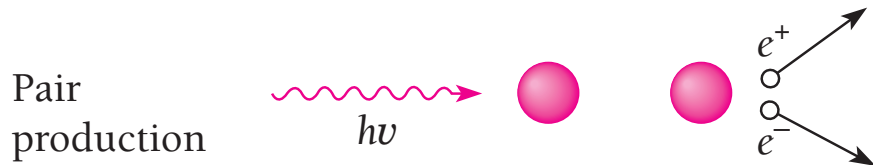
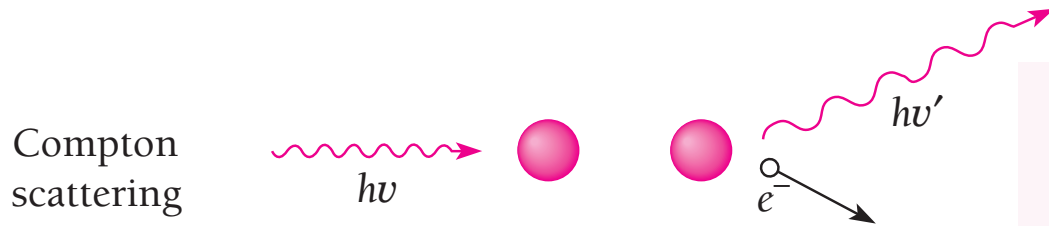
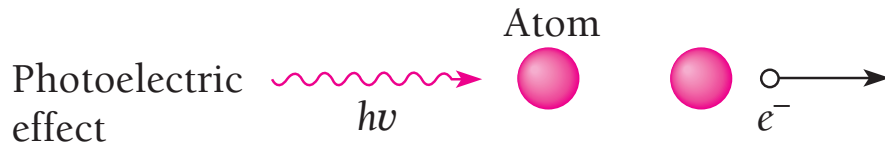
In other words, the two photons from positronium annihilation will move in opposite directions, each with energy 0.511 MeV. This is exactly what is observed experimentally.

6.4 Photon absorption



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The three chief ways in which photons of light, x-rays, and gamma rays interact with matter



6.4 Photon absorption



The intensity I of an x- or gamma-ray beam is equal to the rate at which it transports energy per unit cross-sectional area of the beam.

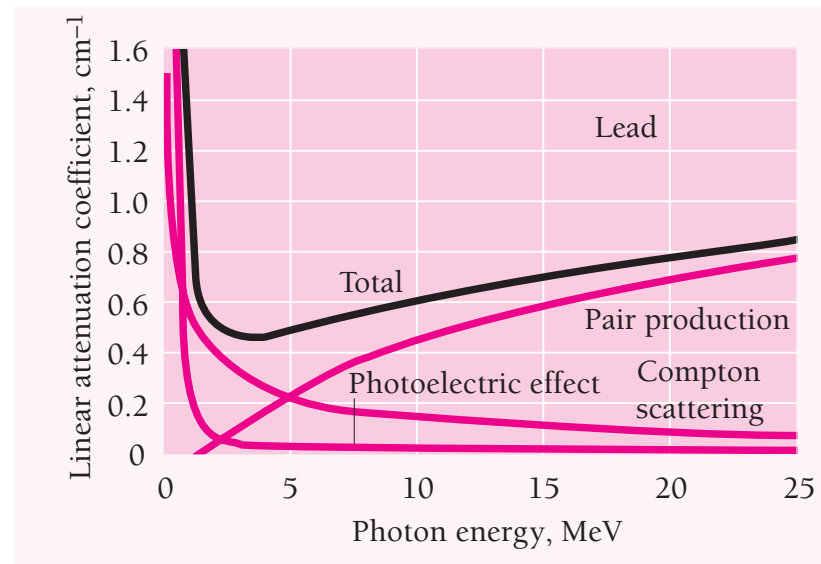
$$-\frac{dI}{I} = \mu dx$$

Radiation intensity I

$$I = I_0 e^{-\mu x}$$

Absorber thickness

$$x = \frac{\ln(I_0/I)}{\mu}$$



Exercise



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1. Which element has a K_{α} x-ray line whose wavelength is 0.180 nm?

Exercise

1. Which element has a K_{α} x-ray line whose wavelength is 0.180 nm?

Solution:

The frequency corresponding to a wavelength of 0.180 nm is

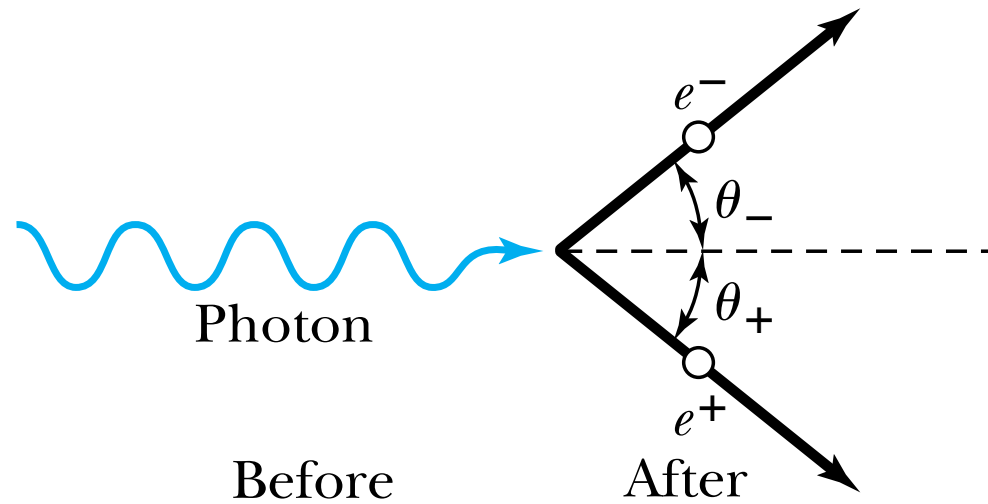
$$\nu = \frac{c}{\lambda} = \frac{3.00 \times 10^8 \text{ m/s}}{1.80 \times 10^{-10} \text{ m}} = 1.67 \times 10^{18} \text{ Hz}$$

Therefore

$$Z - 1 = \sqrt{\frac{4}{3cR}} = \sqrt{\frac{(4)(1.67 \times 10^{18} \text{ Hz})}{(3)(3.00 \times 10^8 \text{ m/s})(1.097 \times 10^7 \text{ m}^{-1})}} = 26$$
$$Z = 27$$

The element with atomic number 27 is cobalt.

2. Show that a photon cannot produce an electron-positron pair in free space as following figure



2. Show that a photon cannot produce an electron-positron pair in free space as following figure

Solution:

Let the total energy and momentum of the electron and the positron be E_- , p_- , and E_+ , p_+ , respectively. The conservation laws are then

$$hf = E_+ + E_-$$

$$\frac{hf}{c} = p_- \cos \theta_- + p_+ \cos \theta_+$$

$$0 = p_- \sin \theta_- - p_+ \sin \theta_+$$



From the second equation, we have

$$hf_{\max} = p_-c + p_+c$$

When we inserted the mass-energy relation to first equation

$$hf = \sqrt{p_+^2c^2 + m^2c^4} + \sqrt{p_-^2c^2 + m^2c^4}$$

Therefore

$$hf > p_-c + p_+c$$

The two frequency equations are inconsistent and cannot simultaneously be valid.

The Physics of Atoms and Quanta

18.1 18.2 18.3 18.4 18.8